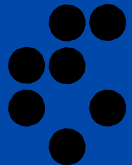


Proton Capture Cross Sections at Low Energy

Matej Lipoglavšek

Jožef Stefan Institute, Ljubljana, Slovenia



Russbach, March 2013

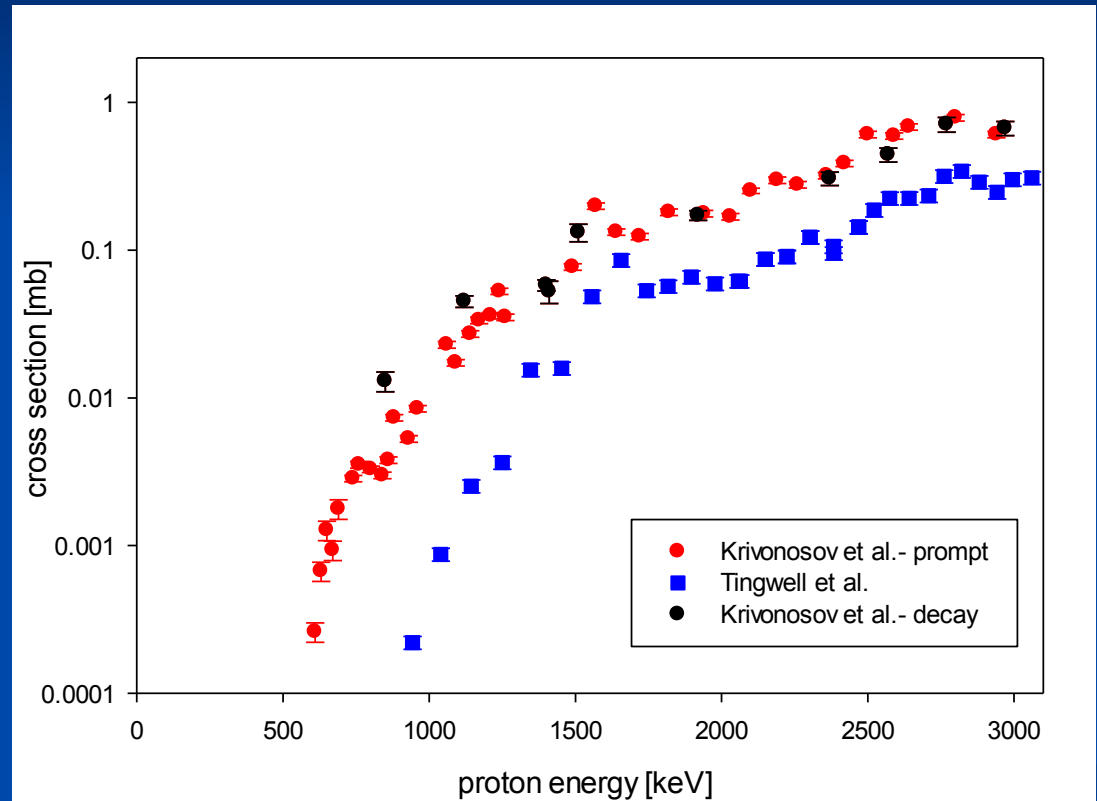
Nuclear Reactions at Low Energies

Due to Coulomb repulsion the cross section σ for charged particle induced nuclear reactions drops rapidly with decreasing beam energy.

$$\sigma(E) = \frac{S(E)}{E} e^{-2\pi\eta},$$

where $\eta = Z_1 Z_2 e^2 / 4\pi\epsilon_0 \hbar \sqrt{2E/\mu}$ is the Sommerfeld parameter. Exponential (Gamow) factor approximates barrier penetration probability.

$^{60}\text{Ni}(p, \gamma)^{61}\text{Cu}$ Cross Section



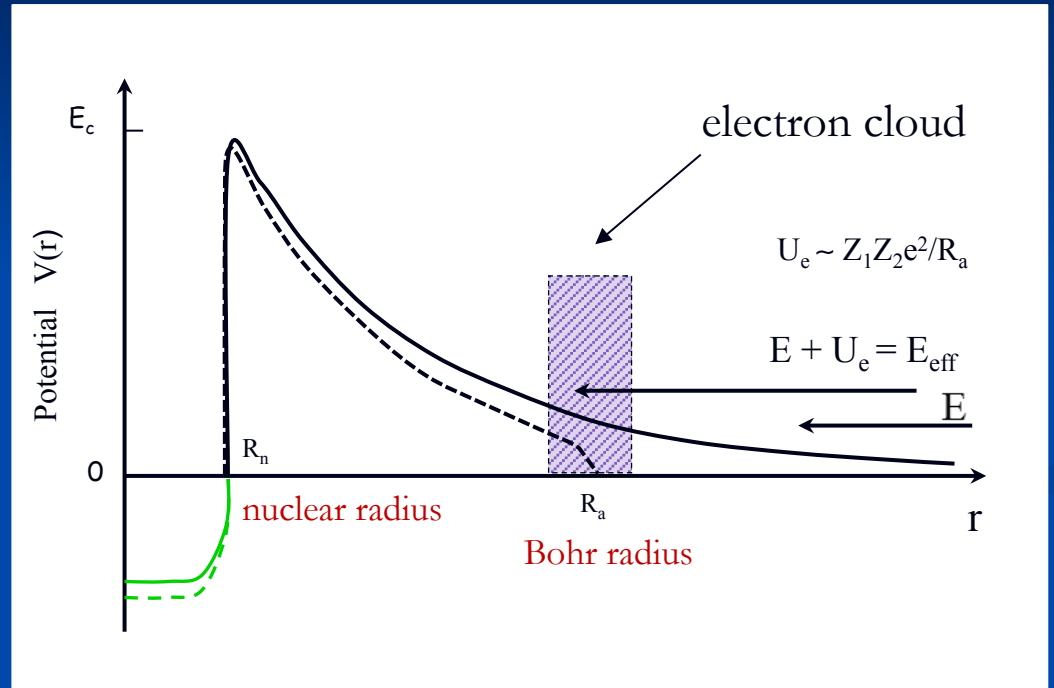
G. A. Krivonosov et al., *Izv. Akad. Nauk SSSR* **41** (1977) 2196.
C. I. W. Tingwell et al., *Nucl. Phys.* **A496** (1989) 127.

Electron Screening

Cross section increases at low energies when the interacting nuclei are not bare. Enhancement factor

$$f(E) = \frac{\sigma(E + U_e)}{\sigma(E)},$$

where U_e is the screening potential.



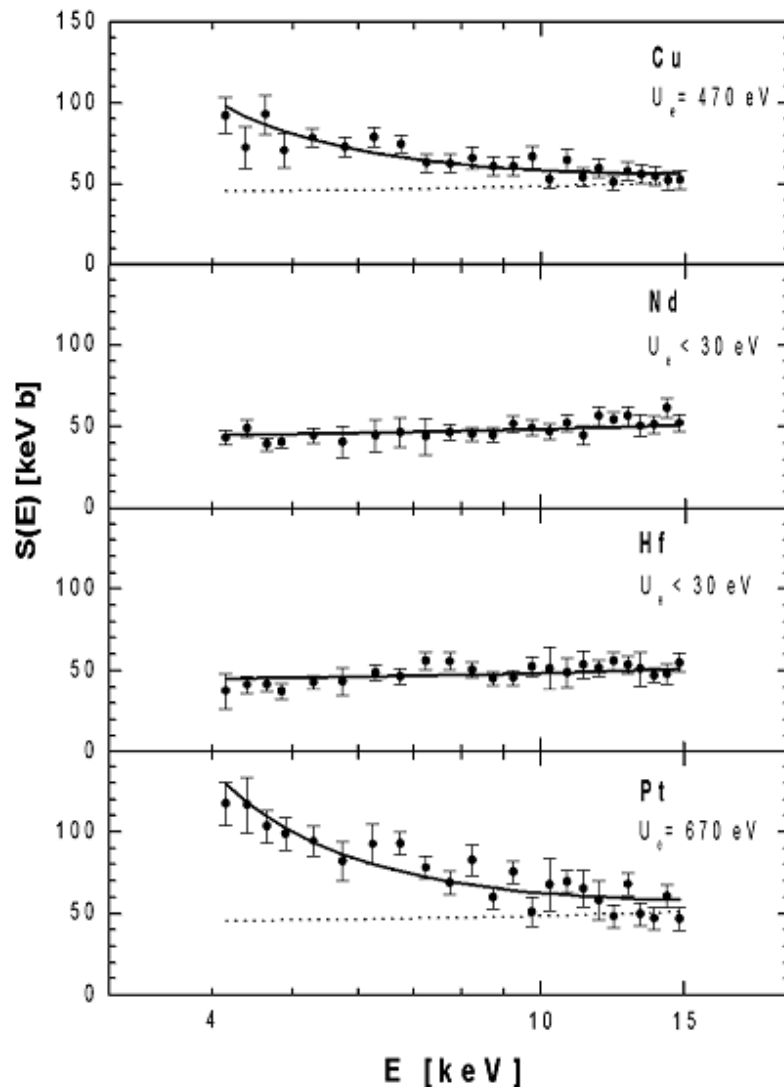
H. J. Assenbaum, K. Langanke and C. Rolfs, Z. Phys. A **327** (1987) 461. 246 citations (Web of Science).

$$\frac{R_n}{R_a} \approx 10^{-5} \Rightarrow U_e = \frac{e^2}{4\pi\epsilon_0 R_a} = 27 \text{ eV for d+d reaction}$$



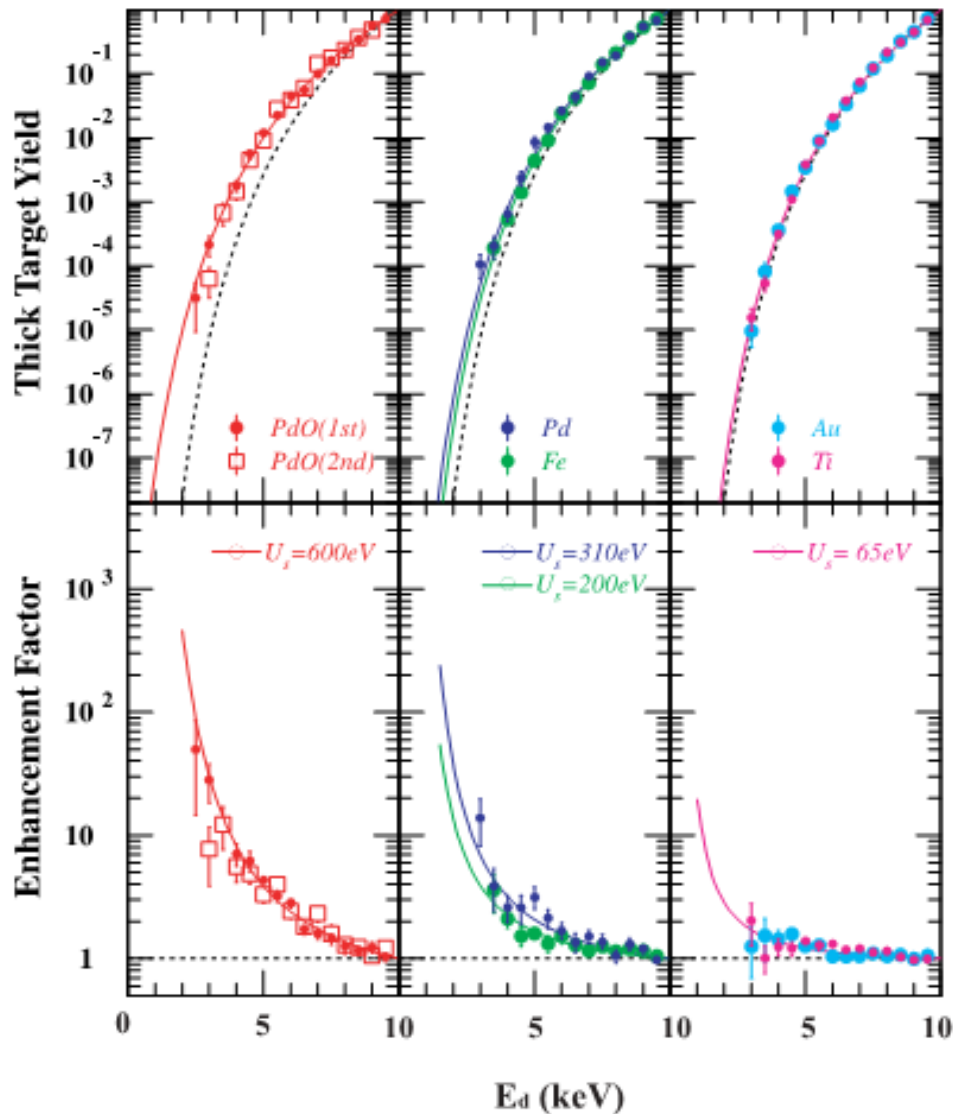
Previous Results 1

for $d(d,p)t$ reaction from F. Raiola et al., Eur. Phys. J. A19 (2004) 283.



Material	U_e (eV) ^(b)	Solubility $1/x$ ^(c)	n_{eff} ^(b)	n_{eff} (Hall) ^(d)
Metals				
Be	180±40	0.08	0.2±0.1	(0.21±0.04)
Mg	440±40	0.11	3.0±0.5	1.8±0.4
Al	520±50	0.26	3.0±0.6	3.1±0.6
V	480±60	0.04	2.1±0.5	(1.1±0.2)
Cr	320±70	0.15	0.8±0.4	(0.20±0.04)
Mn	390±50	0.12	1.2±0.3	(0.8±0.2)
Fe	460±60	0.06	1.7±0.4	(3.0±0.6)
Co	640±70	0.14	3.1±0.7	(1.7±0.3)
Ni	380±40	0.13	1.1±0.2	1.1±0.2
Cu	470±50	0.09	1.8±0.4	1.5±0.3
Zn	480±50	0.13	2.4±0.5	(1.5±0.3)
Sr	210±30	0.27	1.7±0.5	
Nb	470±60	0.13	2.7±0.7	(1.3±0.3)
Mo	420±50	0.12	1.9±0.5	(0.8±0.2)
Ru	215±30	0.18	0.4±0.1	(0.4±0.1)
Rh	230±40	0.09	0.5±0.2	(1.7±0.4)
Pd	800±90	0.03	6.3±1.3	1.1±0.2
Ag	330±40	0.14	1.3±0.3	1.2±0.3
Cd	360±40	0.18	1.9±0.4	(2.5±0.5)
In	520±50	0.02	4.8±0.9	
Sn	130±20	0.08	0.3±0.1	
Sb	720±70	0.13	11±2	
Ba	490±70	0.21	9.9±2.9	
Ta	270±30	0.13	0.9±0.2	(1.1±0.2)
W	250±30	0.29	0.7±0.2	(0.8±0.2)
Re	230±30	0.14	0.5±0.1	(0.3±0.1)
Ir	200±40	0.23	0.4±0.2	(2.2±0.5)
Pt	670±50	0.06	4.6±0.7	3.9±0.8
Au	280±50	0.18	0.9±0.3	1.5±0.3
Tl	550±90	0.01	5.8±1.2	(7.4±1.5)
Pb	480±50	0.04	4.3±0.9	
Bi	540±60	0.12	6.9±1.5	

Previous Results 2



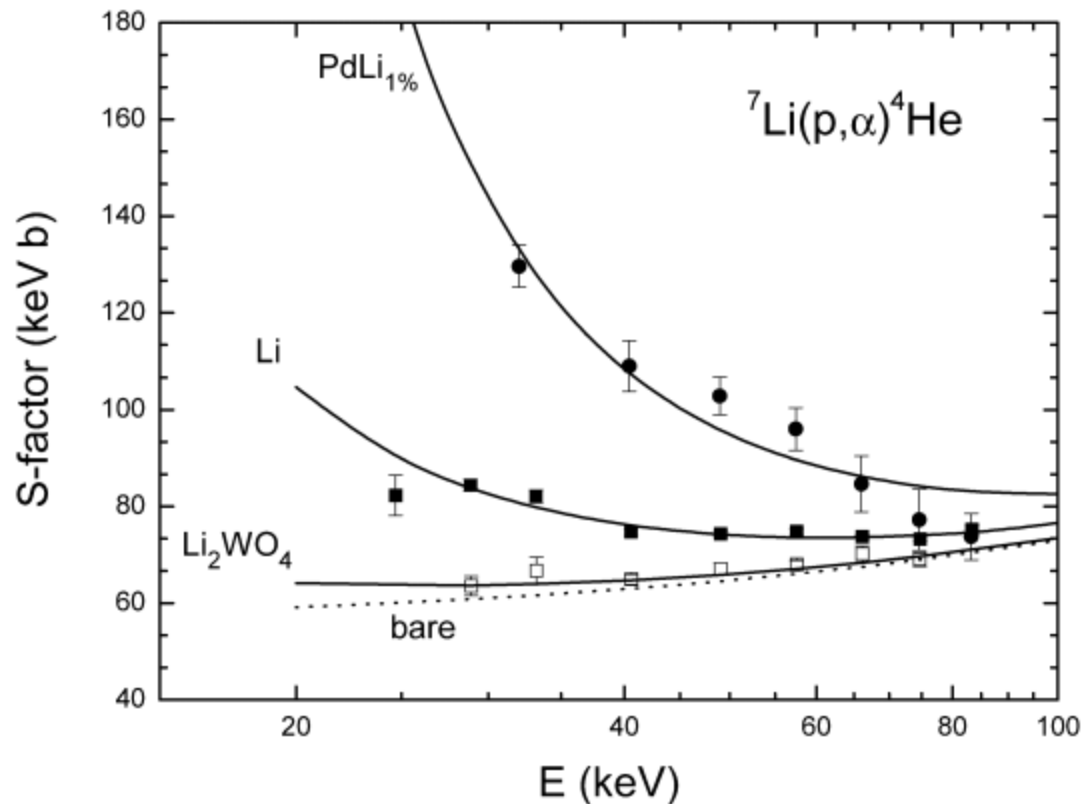
J. Kasagi, Prog. Theo. Phys. Suppl. 154 (2004) 365

for the $d(d,p)t$ reaction
 $U_e = 310 \pm 30$ eV @ 7% H/Pd

\Rightarrow concentration dependence

Previous Results 3

J. Cruz et al., Phys. Lett. B 624 (2005) 181.



For PdLi_{1%}:

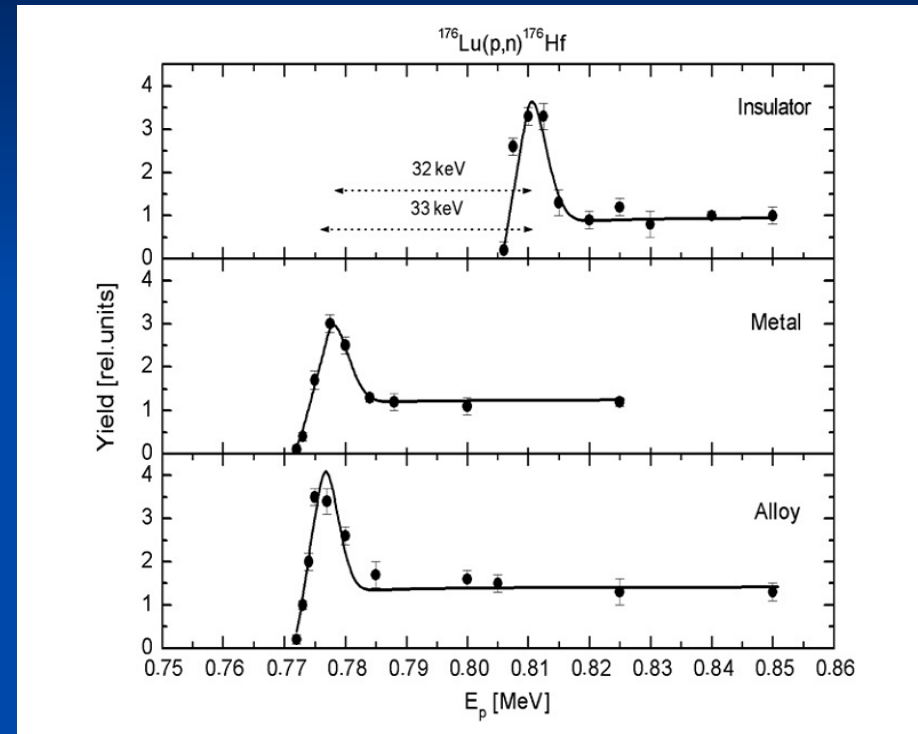
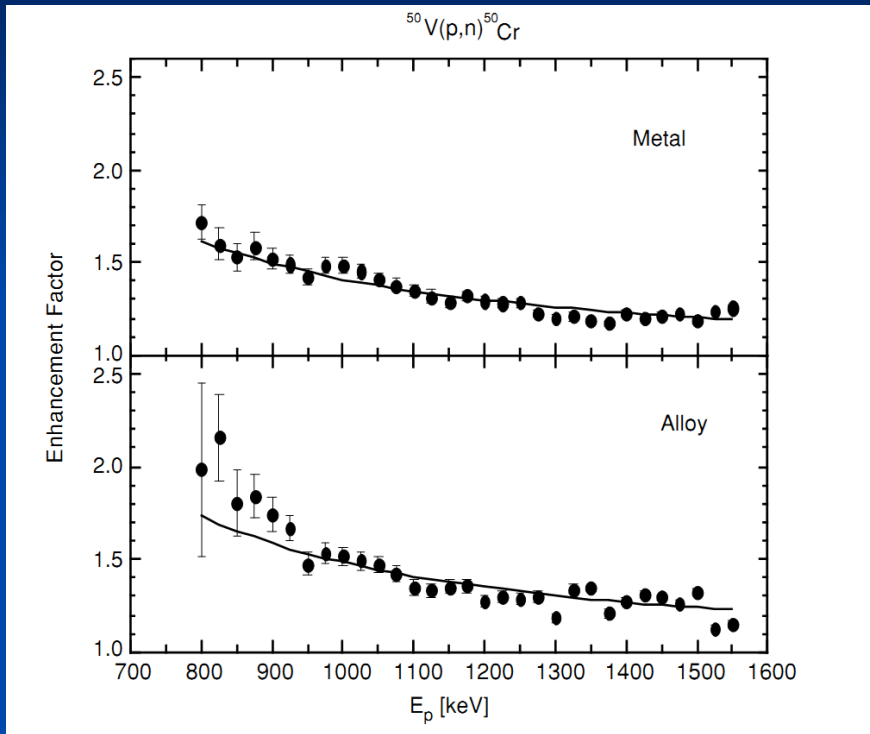
$$S(E) = 0.055 + 0.21E - 0.31E^2$$

[MeV b]

$$U_e = 3.8 \text{ keV}$$

Previous Results 4

K. U. Kettner et al., J. Phys. G **32** (2006) 489.



$^{50}\text{V}(p,n)^{50}\text{Cr}$ reaction in different environments: VO_2 **insulator**, V metal and $\text{PdV}_{10\%}$ alloy. Relative to the insulator, metal and alloy showed a large screening potential of $U_e = 27$ and **34 keV**.

$^{176}\text{Lu}(p,n)^{176}\text{Hf}$ reaction in Lu_2O_3 **insulator**, Lu metal and $\text{PdLu}_{10\%}$ alloy; there is a narrow resonance and a shift in proton resonance energy of $U_e = 32$ and **33 keV** for the metal and alloy, respectively, relative to the **insulator**.

$$\Rightarrow U \downarrow e \propto Z$$

Jožef Stefan Institute

10th Russbach School

on Nuclear Astrophysics

March 10 – 16, 2013

Russbach, Austria

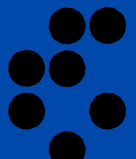
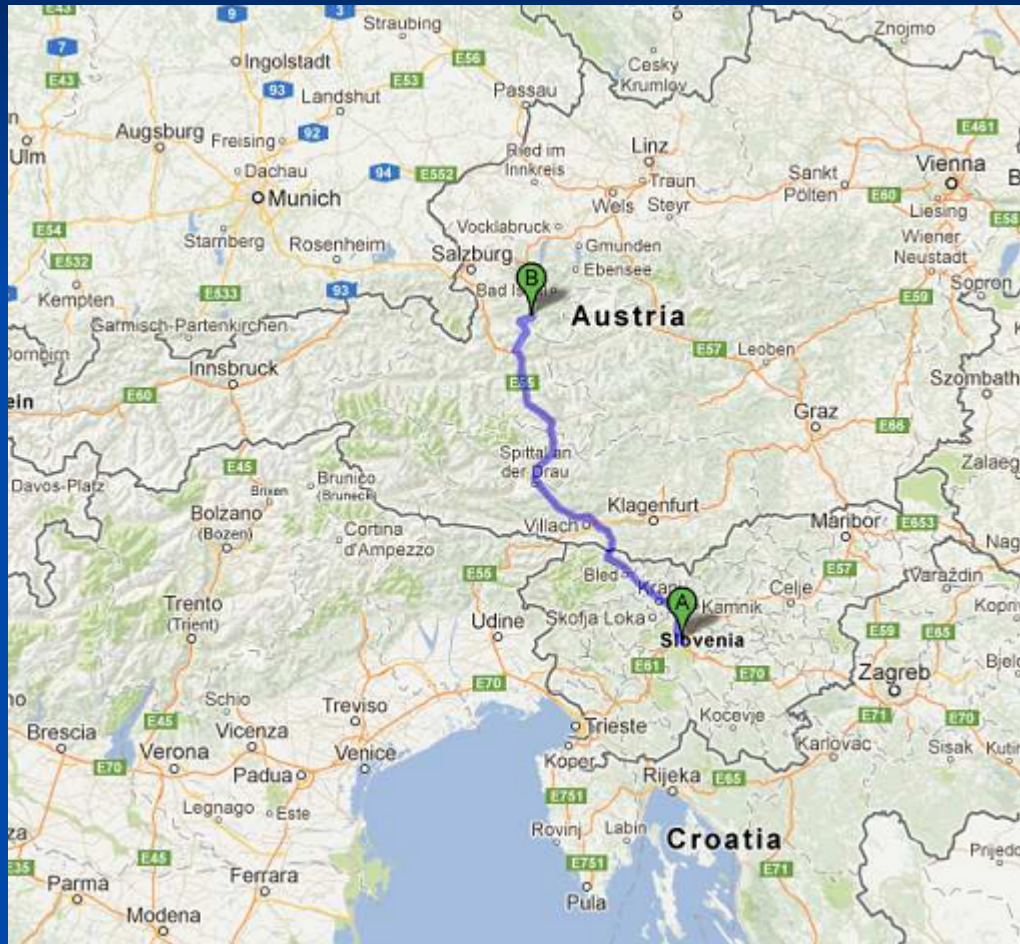


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300 km

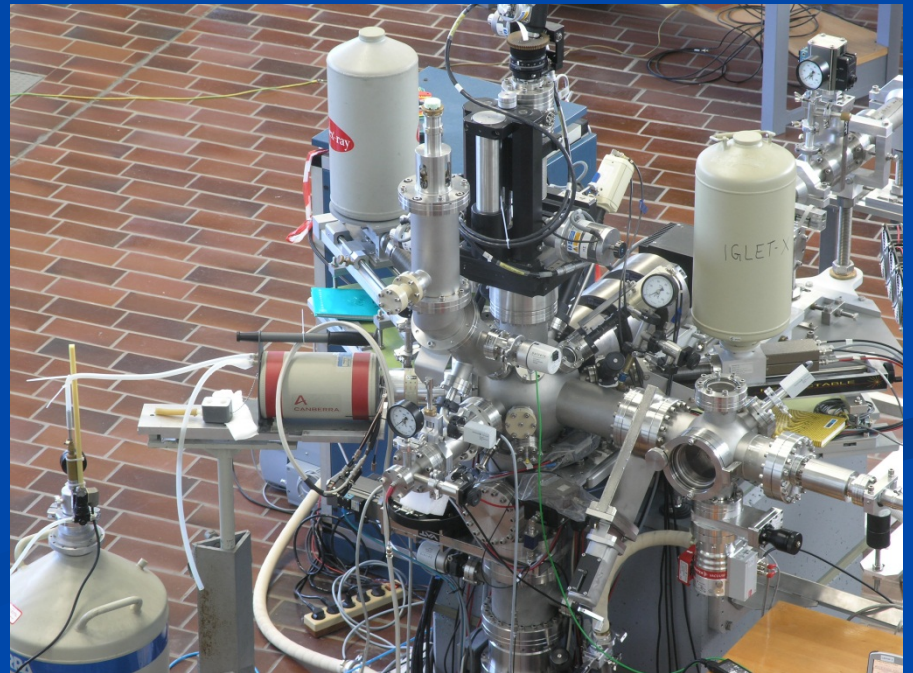


Measurements @ JSI



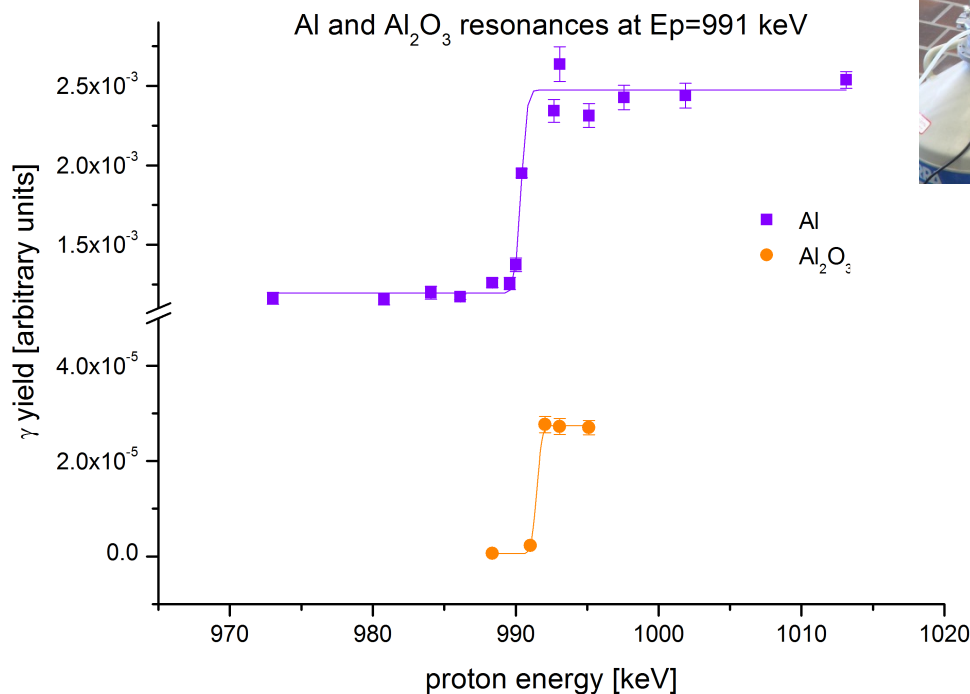
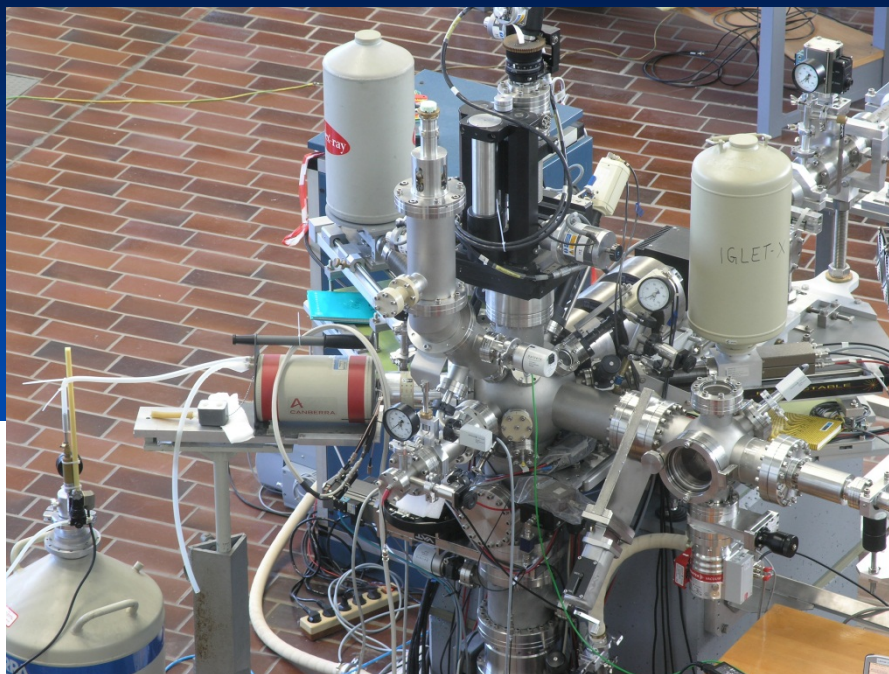
2 MV Tandem van de Graaf accelerator

Experimental setup:
2 X-ray detectors,
1 γ -ray detector,
charged particle detectors,
neutron detector.



Electron screening in aluminum

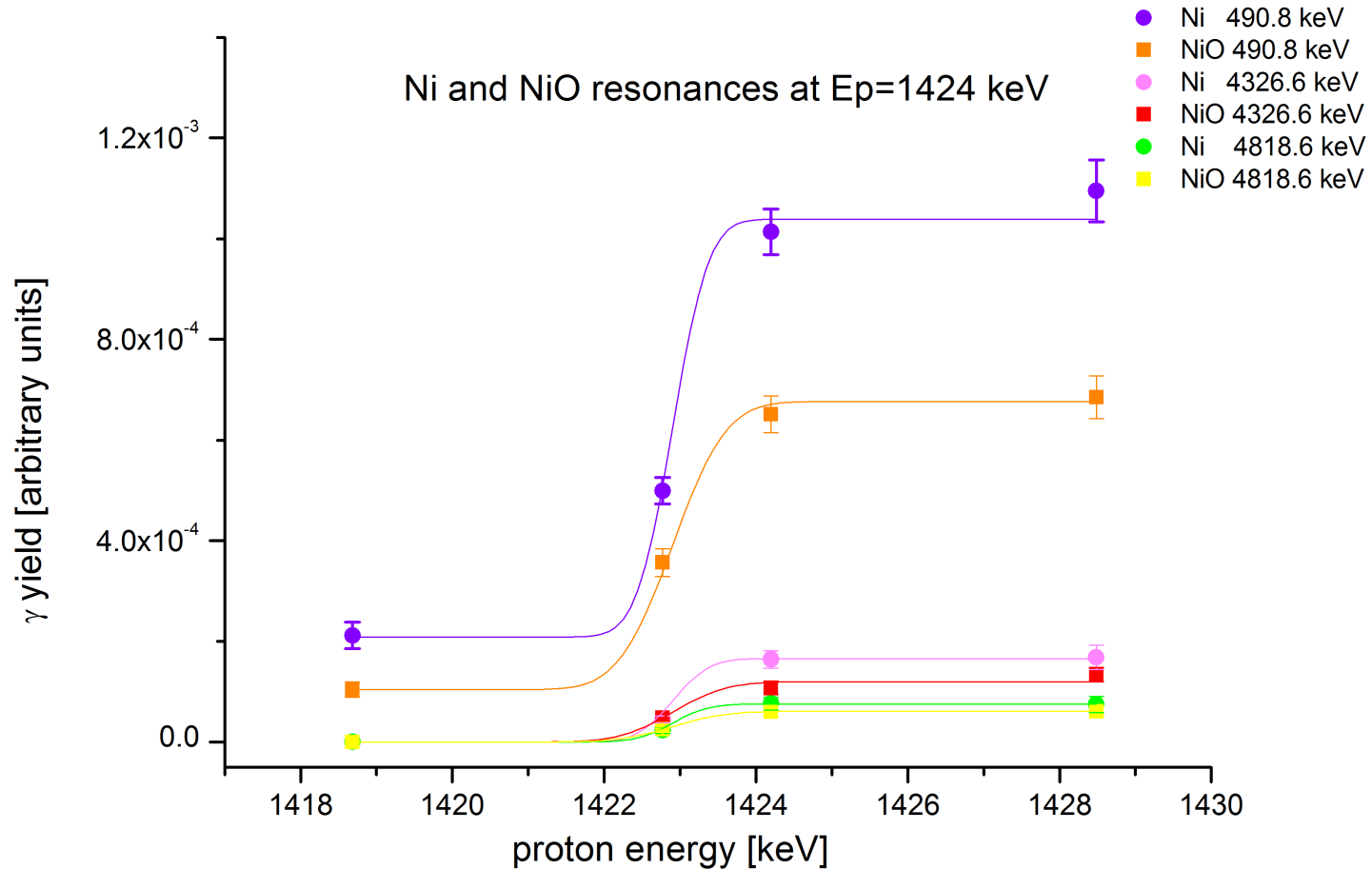
(p,γ) , $(p,p'\gamma)$ and $(p,n\gamma)$ reactions were studied on natural Ni and Al metals and NiO and Al_2O_3 insulators.



Characteristic γ rays were measured by a Ge detector:
1779 keV γ ray from $^{27}\text{Al}(p,\gamma)^{28}\text{Si}$ reaction

Nickel

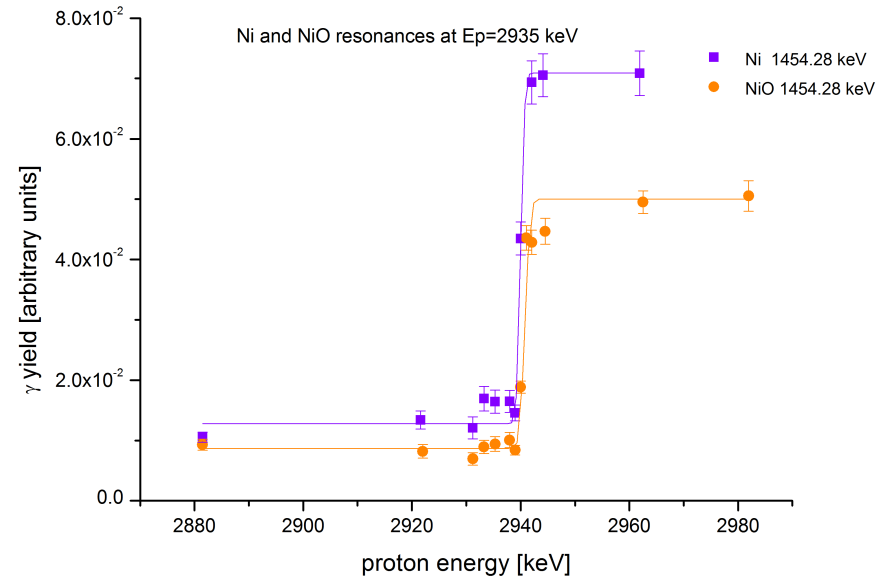
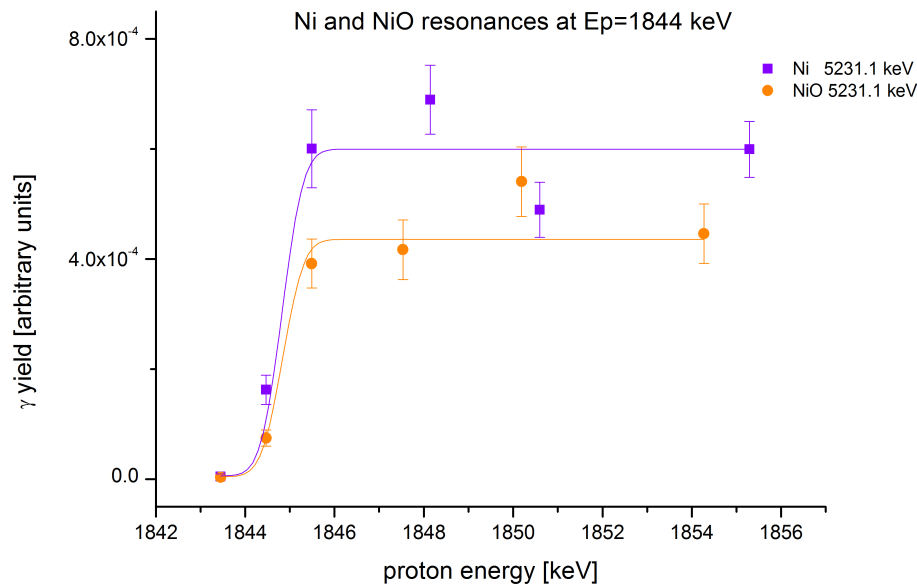
3 γ rays from $^{58}\text{Ni}(p,\gamma)^{59}\text{Cu}$ reaction



Nickel

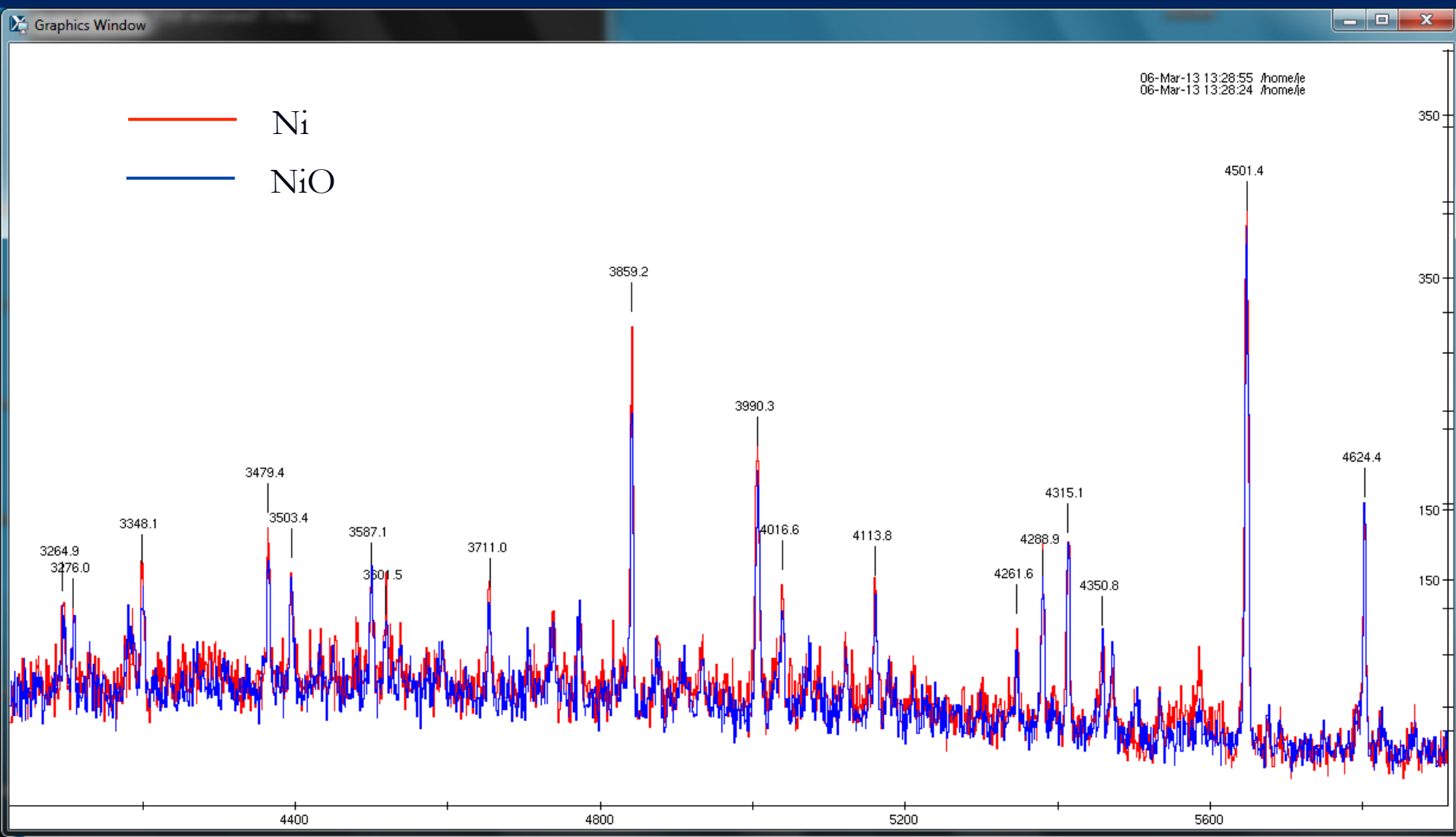
1454 keV γ ray from $^{58}\text{Ni}(p,p'\gamma)^{58}\text{Ni}$ reaction

5231 keV γ ray from $^{58}\text{Ni}(p,\gamma)^{59}\text{Cu}$ reaction

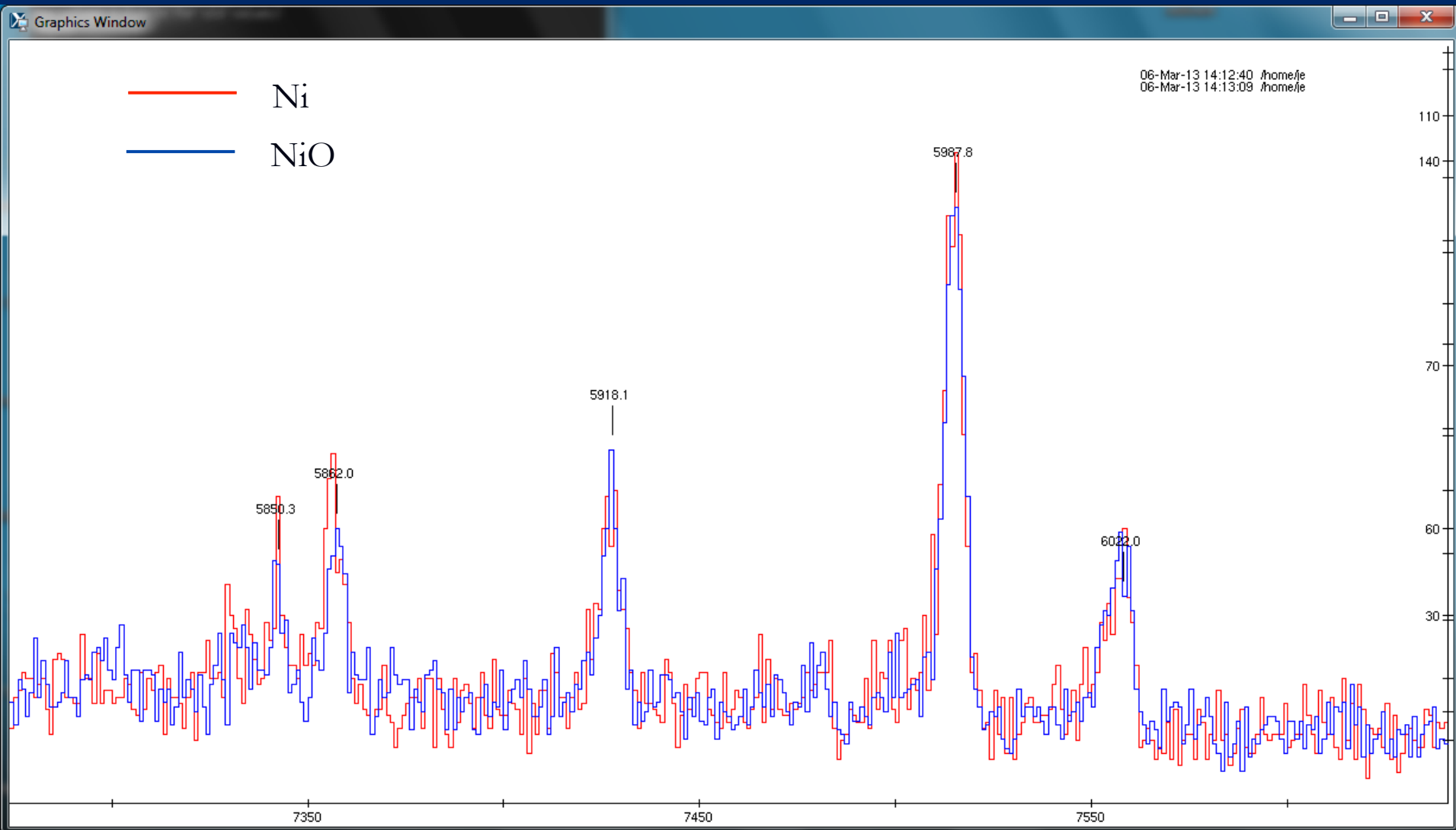


- No shift in resonance energy
- Difference in resonance strength

$^{58}\text{Ni}(p, \gamma)^{59}\text{Cu}$



$^{60}\text{Ni}(p, \gamma)^{61}\text{Cu}$

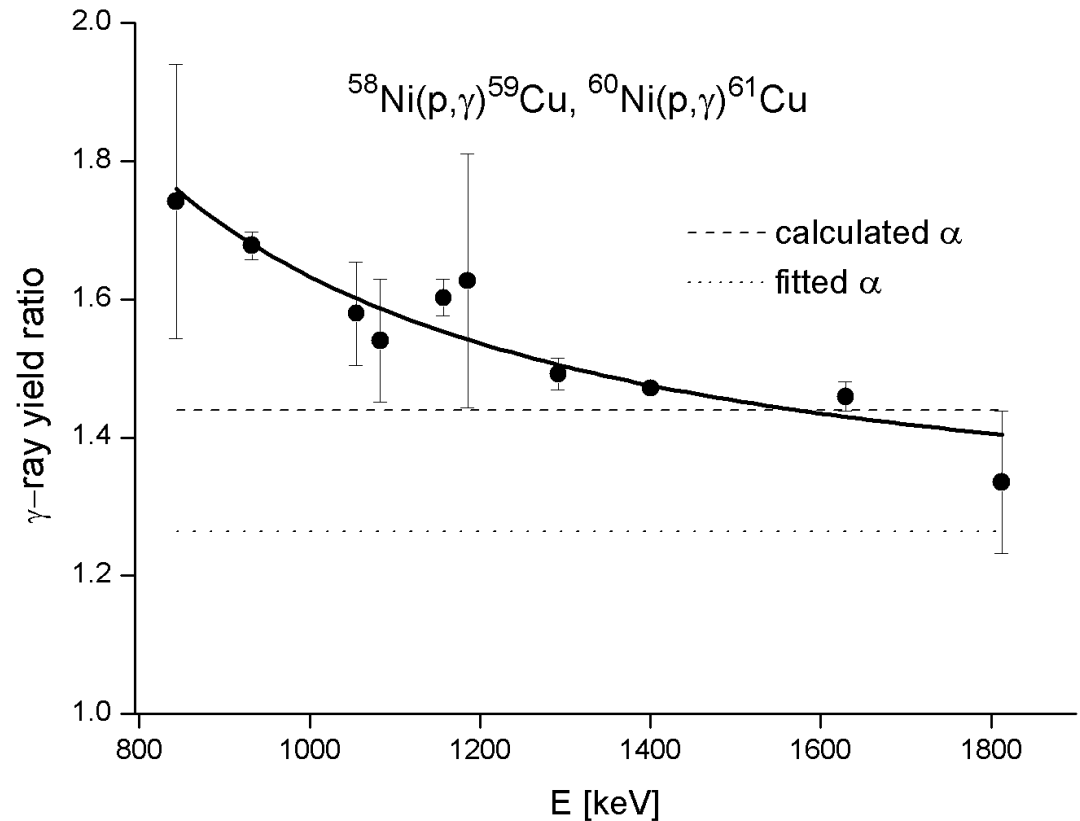


Enhancement factor

$$f(E) = \frac{\sigma(E + U_e)}{\sigma(E)}$$

$$f(E) = \alpha \frac{N_\gamma / N_p(\text{Ni})}{N_\gamma / N_p(\text{NiO})}$$

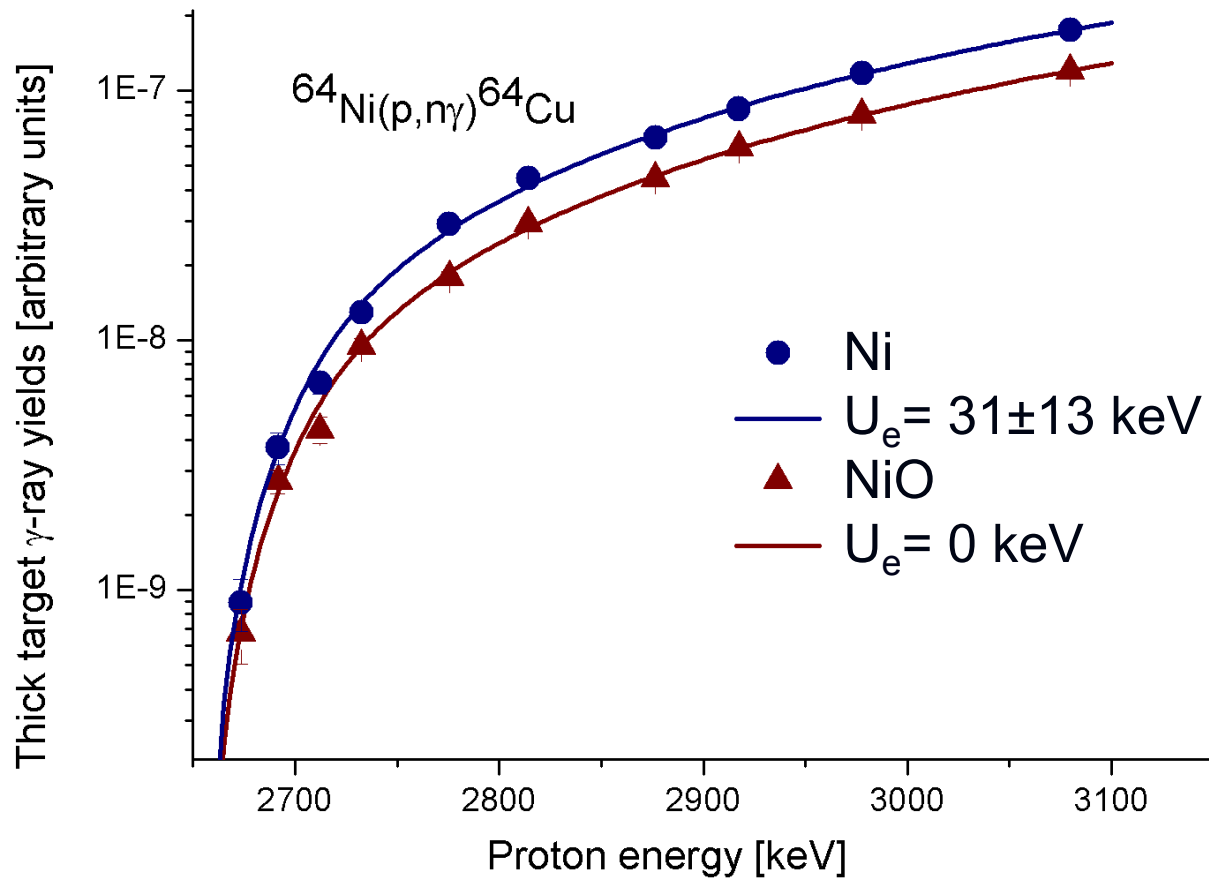
$$\alpha = \frac{M_{\text{NiO}} \frac{dE}{dx\rho}(\text{NiO})}{M_{\text{Ni}} \frac{dE}{dx\rho}(\text{Ni})}$$



$$U_e = 19 \pm 4 \text{ keV}; \alpha = 1.26 \pm 0.02$$

$^{64}\text{Ni}(p,n)^{64}\text{Cu}$ reaction

$E_\gamma = 159 \text{ keV}$

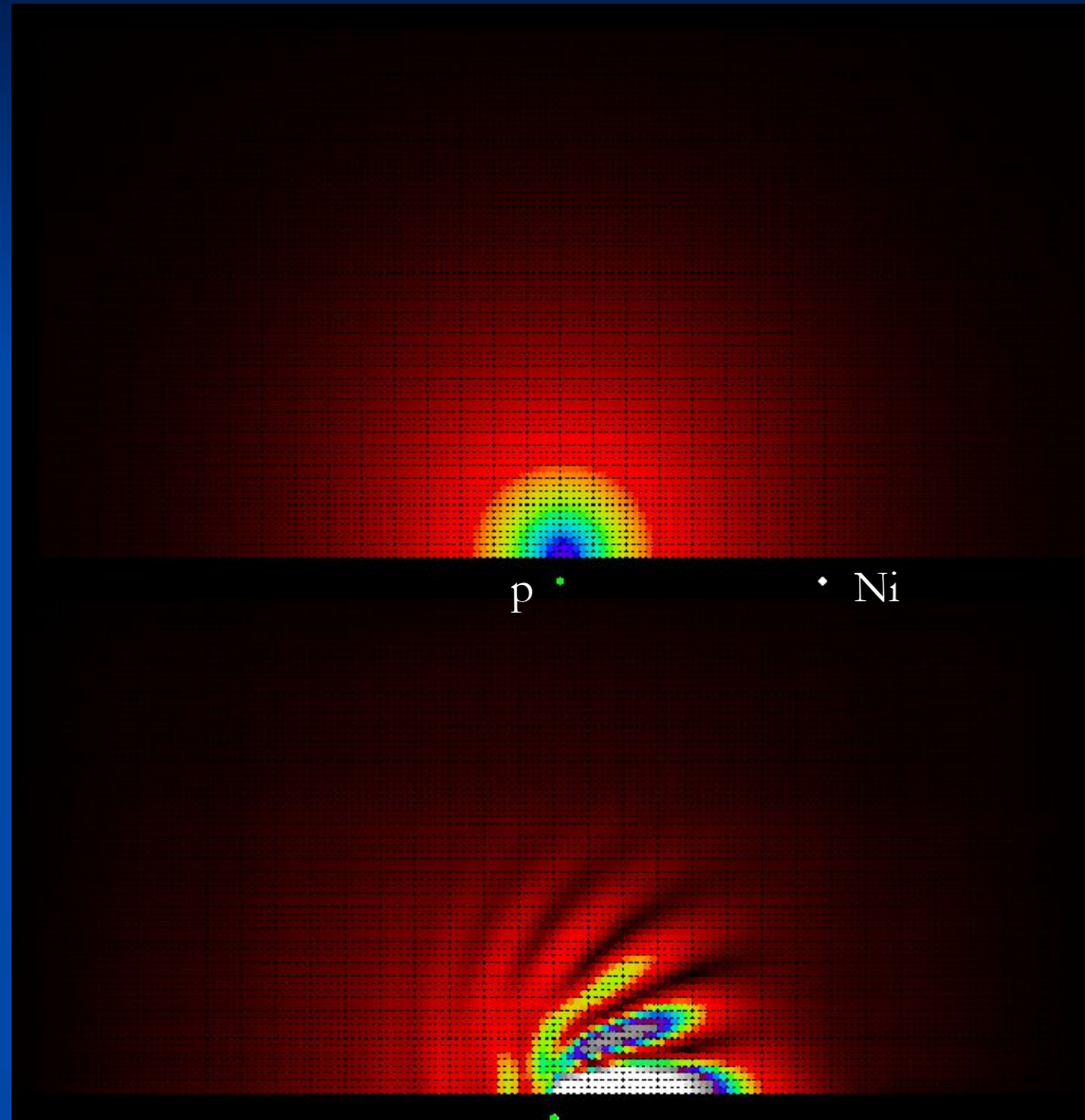


Electron screening calculation

Solving time dependent 3D
Schrödinger equation on a lattice
for the hydrogen electron in collision
of neutral hydrogen and nickel
atoms at hydrogen energy of
1.6 MeV.

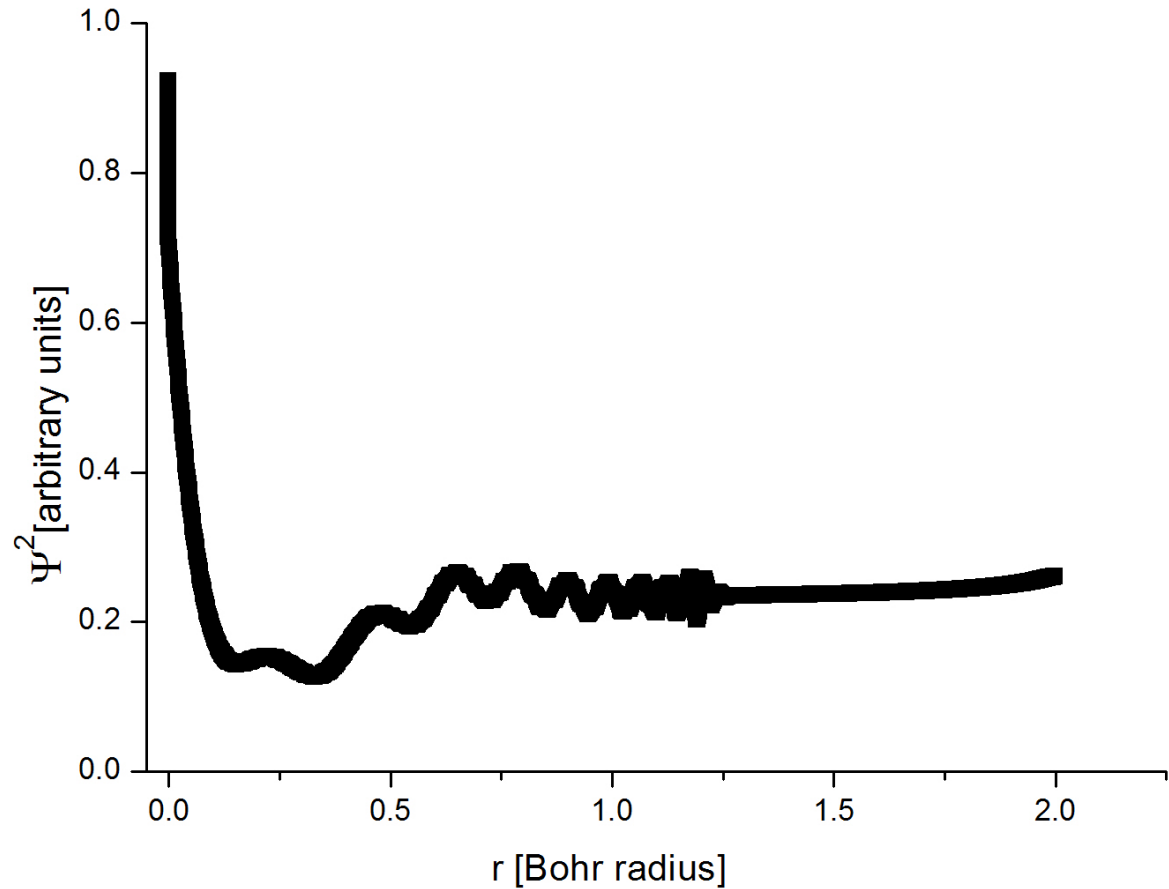
$t=0$

$t=\text{collision}$



Electron density

Electron density at the place of the proton as a function of the distance between the proton and the nickel nucleus.

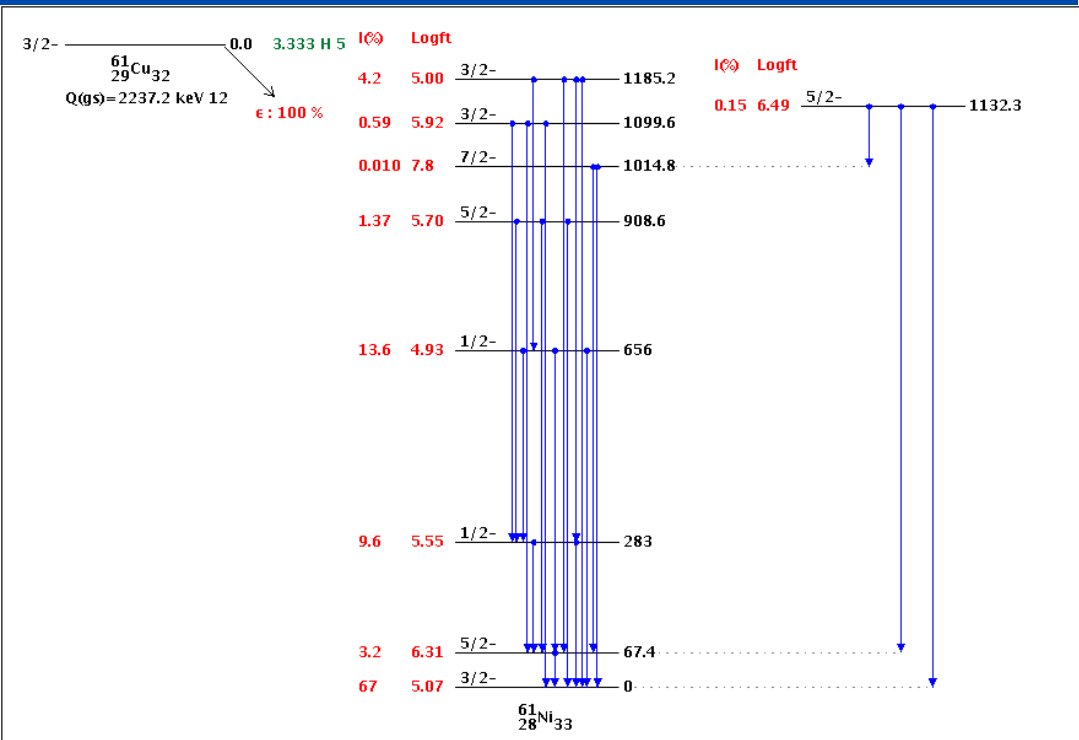
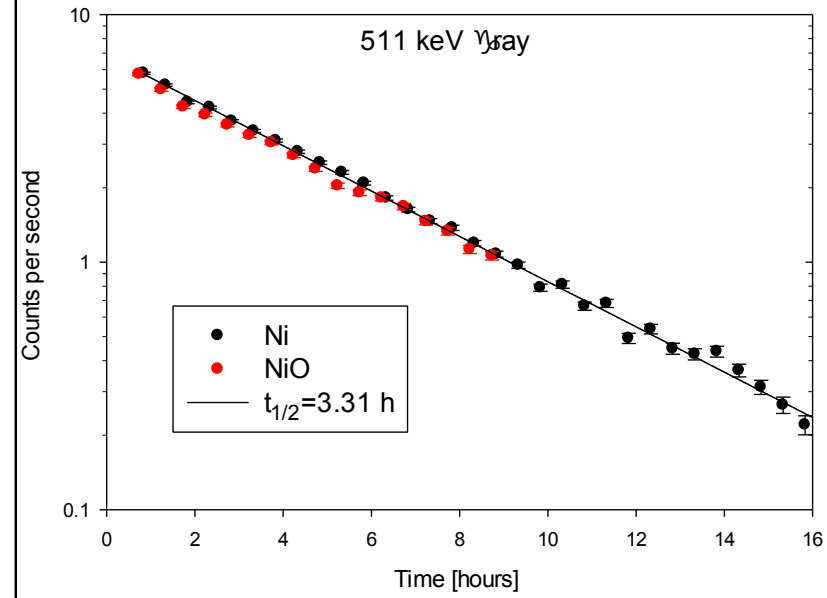
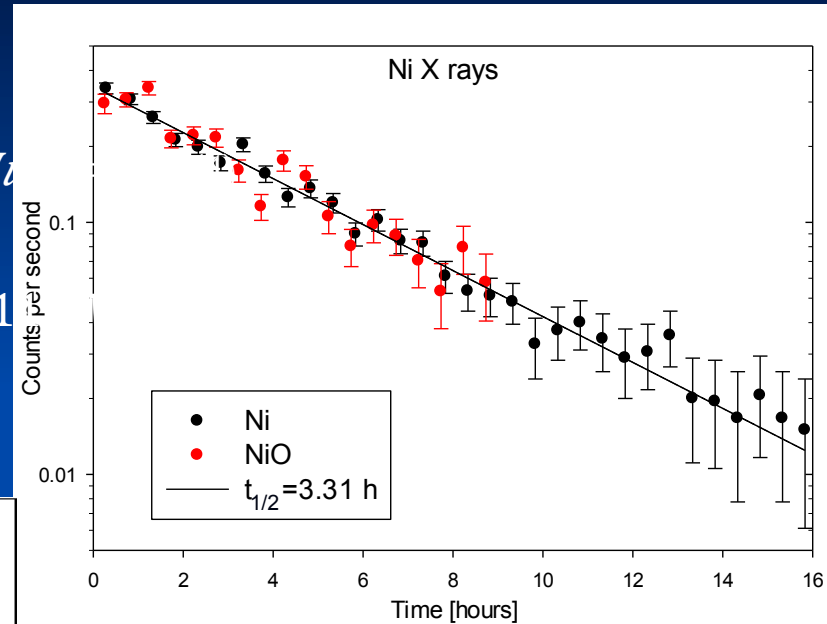


β^+ /EC Decay of ^{61}Cu

^{61}Cu source	Half-life [h]
Ni metal	3.31(1)
NiO insulator	3.34(2)
Nucl. Data Sheets	3.333(5)

$N_{\downarrow X}(\text{Ni})/N_{\downarrow X}(\text{NiO})$

$N_{\downarrow 511}(\text{Ni})/N_{\downarrow 511}(\text{NiO})$



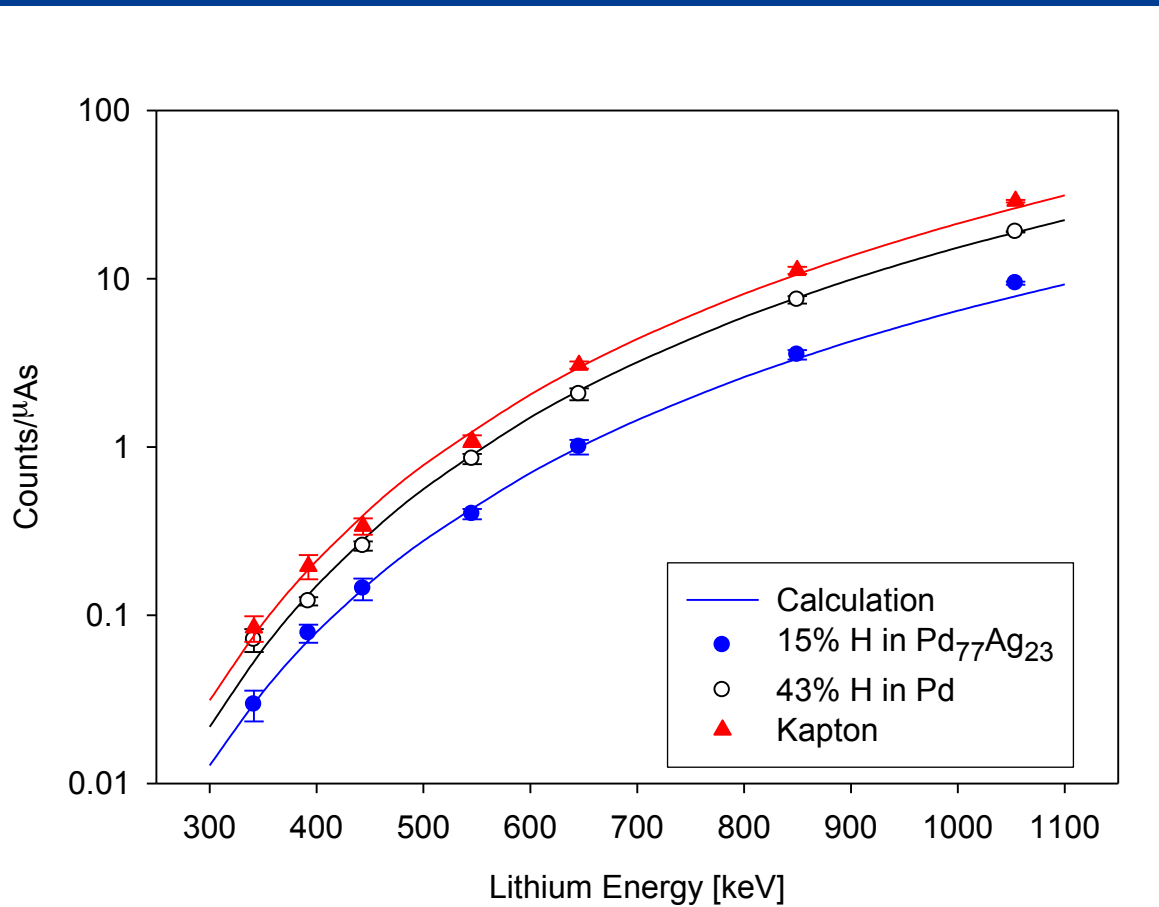
Conclusions

- All measured cross sections in metallic targets are not bare, but include electron screening. The same is true for resonance strengths.
- $U_e(Z=1;d,p)=200-800$ eV, $U_e(Z=3;p,\alpha)=3.8 \pm 0.5$ keV, $U_e(Z=23;p,n)=27 \pm 9$ keV, $U_e(Z=28;p,n)=31 \pm 13$ keV, $U_e(Z=28;p,\gamma)=19 \pm 4$ keV, $U_e(Z=72;p,n)=32 \pm 2$ keV.
- Electron screening is not a static but rather a dynamic effect, so the parameterization with a screening potential is questionable.
- Electron screening **IS** important in plasma.

Thanks to: Jelena Gajević, Toni Petrovič, Miha Škof, Andrej Likar, Žiga Šmit, Matjaž Vencelj, Primož Pelicon, Primož Vavpetič, Drago Brodnik, Aleksandra Cvetinović, Alberto Sanchez Ortiz

Thick Target Yields

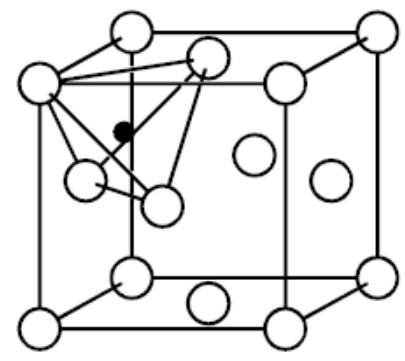
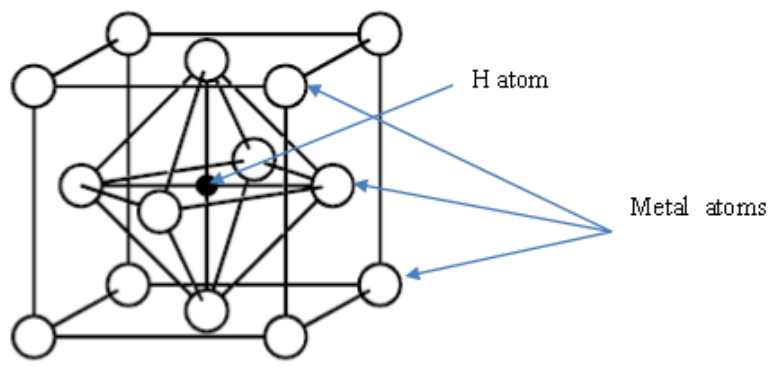
$$\alpha\text{-particle yield calculation: } N_{\alpha} = 2N_{Li} \frac{\rho N_A}{M} \int_{E_0}^0 \Omega W \frac{\sigma(E)}{dE_{Li}/dx} dE_{Li}$$



dE_{Li}/dx stopping power
 Ω efficiency
 W ang. distribution

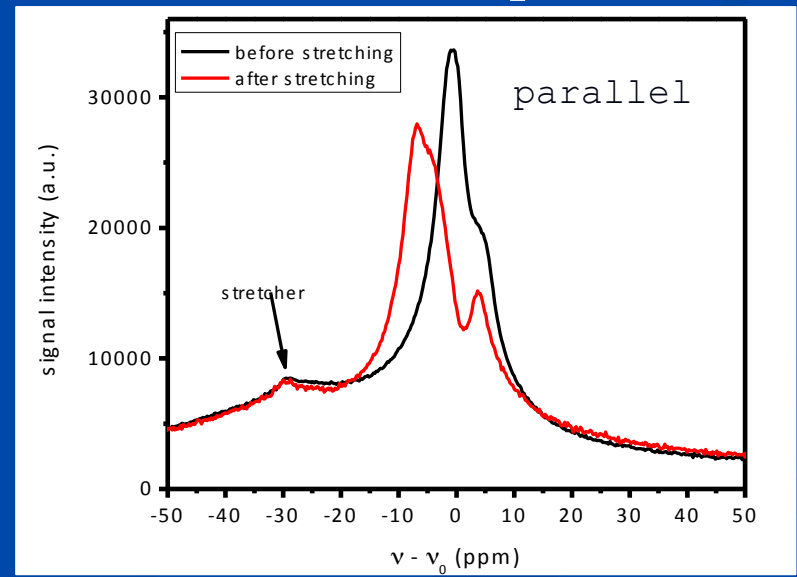
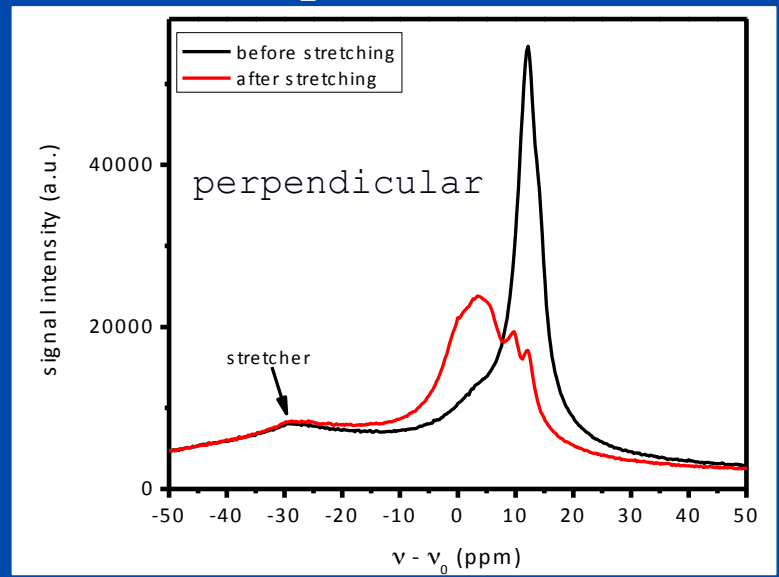
Reaction: ${}^1\text{H}({}^7\text{Li}, \alpha){}^4\text{He}$

Electron density



face
centered
cubic
lattice

Octahedral position $\xrightarrow{\text{stretching}}$ Tetrahedral position



^1H NMR lineshapes measured by Hahn echo at $\nu_0 = 100$ MHz of a $47 \mu\text{m}$ thick Pd foil.